Do geometric invariants of preferred extremals define topological invariants of space-time surface and code for quantum physics?

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Abstract

The recent progress in the understanding of preferred extremals of Kähler action leads to the conclusion that they satisfy Einstein-Maxwell equations with cosmological term with Newton's constant and cosmological constant predicted to have a spectrum. One particular implication is that preferred extremals have a constant value of Ricci scalar. The implications of this are expected to be very powerful since it is known that $D > 2$-dimensional manifolds allow a constant curvature metric with volume and other geometric invariants serving as topological invariants. Also the the possibly discrete generalization of Ricci flow playing key role in manifold topology to Maxwell flow is very natural, and the connections with the geometric description of dissipation, self-organization, transition to chaos and also with coupling constant evolution are highly suggestive. A further fascinating possibility inspired by quantum classical correspondence is quantum ergodicity (QE): the statistical geometric properties of preferred extremals code for various correlations functions of zero energy states defined as their superpositions so that any preferred extremal in the superposition would serve as a representative of the zero energy state. QE would make possible to deduce correlation functions and S-matrix from the properties of single preferred extremal.

Contents

1 Introduction 1

2 Preferred extremals of Kähler action as manifolds with constant Ricci scalar whose geometric invariants are topological invariants 2

3 Is there a connection between preferred extremals and AdS/CFT correspondence? 3

4 Generalizing Ricci flow to Maxwell flow for 4-geometries and Kähler flow for space-time surfaces 5

4.1 Ricci flow and Maxwell flow for 4-geometries 5

4.2 Maxwell flow for space-time surfaces 7

4.3 Dissipation, self organization, transition to chaos, and coupling constant evolution 8

4.4 Does a 4-D counterpart of thermodynamics make sense? 9

5 Could correlation functions, S-matrix, and coupling constant evolution be coded the statistical properties of preferred extremals? 10

1 Introduction

The recent progress in the understanding of preferred extremals [K2, K5] led to a reduction of the field equations to conditions stating for Euclidian signature the existence of Kähler metric. The
resulting conditions are a direct generalization of corresponding conditions emerging for the string world sheet and stating that the 2-metric has only non-diagonal components in complex/hypercomplex coordinates. Also energy momentum of Kähler action and has this characteristic (1,1) tensor structure. In Minkowskian signature one obtains the analog of 4-D complex structure combining hyper-complex structure and 2-D complex structure.

The construction lead also to the understanding of how Einstein’s equations with cosmological term follow as a consistency condition guaranteeing that the covariant divergence of the Maxwell’s energy momentum tensor assignable to Kähler action vanishes. This gives $T = kG + \Lambda g$. By taking trace a further condition follows from the vanishing trace of $T$:

$$R = \frac{4\Lambda}{k}.$$  

That any preferred extremal should have a constant Ricci scalar proportional to cosmological constant is very strong prediction. Note that the accelerating expansion of the Universe would support positive value of $\Lambda$. Note however that both $\Lambda$ and $k \propto 1/G$ are both parameters characterizing one particular preferred extremal. One could of course argue that the dynamics allowing only constant curvature space-times is too simple. The point is however that particle can topologically condense on several space-time sheets meaning effective superposition of various classical fields defined by induced metric and spinor connection.

The following considerations demonstrate that preferred extremals might be seen as canonical representatives for the constant curvature manifolds playing central role in Thurston’s geometrization theorem known also as hyperbolization theorem. The generalization of the notion of Ricci flow to Maxwell flow in the space of metrics and further to Kähler flow for preferred extremals in turn gives a rather detailed vision about how preferred extremals organize to one-parameter orbits. It is quite possible that Kähler flow is actually discrete. The natural interpretation is in terms of dissipation and self organization.

Quantum classical correspondence suggests that this line of thought could be continued even further: could the geometric invariants of the preferred extremals could code not only for space-time topology but also for quantum physics? How to calculate the correlation functions and coupling constant evolution has remained a basic unresolved challenge of quantum TGD. Could the correlation functions be reduced to statistical geometric invariants of preferred extremals? The latest (means the end of 2012) and perhaps the most powerful idea hitherto about coupling constant evolution is quantum classical correspondence in statistical sense stating that the statistical properties of a preferred extremal in quantum superposition of them are same as those of the zero energy state in question. This principle would be quantum generalization of ergodic theorem stating that the time evolution of a single member of ensemble represents the ensemble statistically. This principle would allow to deduce correlation functions and S-matrix from the statistical properties of single preferred extremal alone using classical intuition. Also coupling constant evolution would be coded by the statistical properties of the representative preferred extremal.
the parameter $R = 4\Lambda/k$ and also $\Lambda$ and $k$ separately characterize the equivalence class of preferred extremals as is also physically clear.

Also the volume of the space-time sheet closed inside causal diamond $CD$ remains constant along the orbits of the flow and thus characterizes the space-time surface. $\Lambda$ and even $k \propto 1/G$ can indeed depend on space-time sheet and p-adic length scale hypothesis suggests a discrete spectrum for $\Lambda/k$ expressible in terms of p-adic length scales: $\Lambda/k \propto 1/L^2_p$ with $p \simeq 2^k$ favored by p-adic length scale hypothesis. During cosmic evolution the p-adic length scale would increase gradually. This would resolve the problem posed by cosmological constant in GRT based theories.

2. One could also see the preferred extremals as 4-D counterparts of constant curvature 3-manifolds in the topology of 3-manifolds. An interesting possibility raised by the observed negative value of $\Lambda$ is that most 4-surfaces are constant negative curvature 4-manifolds. By a general theorem of Mostow rigidity theorem [A3] finite-volume hyperbolic manifold is unique for $D > 2$ and determined by the fundamental group of the manifold. Since the orbits under the Kähler flow preserve the curvature scalar the manifolds at the orbit must represent different imbeddings of one and hyperbolic 4-manifold. In 2-D case the moduli space for hyperbolic metric for a given genus $g > 0$ is defined by Teichmüller parameters and has dimension $6(g-1)$. Obviously the exceptional character of $D = 2$ case relates to conformal invariance. Note that the moduli space in question plays a key role in p-adic mass calculations [K3].

In the recent case Mostow rigidity theorem could hold true for the Euclidian regions and maybe generalize also to Minkowskian regions. If so then both ”topological” and ”geometro” in ”Topological GeometroDynamics” would be fully justified. The fact that geometric invariants become topological invariants also conforms with ”TGD as almost topological QFT” and allows the notion of scale to find its place in topology. Also the dream about exact solvability of the theory would be realized in rather convincing manner.

These conjectures are the main result independent of whether the generalization of the Ricci flow discussed in the sequel exists as a continuous flow or possibly discrete sequence of iterates in the space of preferred extremals of Kähler action. My sincere hope is that the reader could grasp how far reaching these result really are.

3 Is there a connection between preferred extremals and AdS$_4$/CFT correspondence?

The preferred extremals satisfy Einstein Maxwell equations with a cosmological constant and have negative scalar curvature for negative value of $\Lambda$. 4-D space-times with hyperbolic metric provide canonical representation for a large class of four-manifolds and an interesting question is whether these spaces are obtained as preferred extremals and/or vacuum extremals.

4-D hyperbolic space with Minkowski signature is locally isometric with AdS$_4$. This suggests at connection with AdS$_4$/CFT correspondence of M-theory. The boundary of AdS would be now replaced with 3-D light-like orbit of partonic 2-surface at which the signature of the induced metric changes. The metric 2-dimensionality of the light-like surface makes possible generalization of 2-D conformal invariance with the light-like coordinate taking the role of complex coordinate at light-like boundary. AdS could represent a special case of a more general family of space-time surfaces with constant Ricci scalar satisfying Einstein-Maxwell equations and generalizing the AdS$_4$/CFT correspondence. There is however a strong objection from cosmology: the accelerated expansion of the Universe requires positive value of $\Lambda$ and favors De Sitter Space $dS_4$ instead of AdS$_4$.

These observations provide motivations for finding whether AdS$_4$ and/or $dS_4$ allows an imbedding as a vacuum extremal to $M^4 \times S^2 \subset M^4 \times CP_2$, where $S^2$ is a homologically trivial geodesic sphere
3. Is there a connection between preferred extremals and AdS$_4$/CFT correspondence?  

of CP$_2$. It is easy to guess the general form of the imbedding by writing the line elements of, $M^4$, $S^2$, and AdS$_4$.  

1. The line element of $M^4$ in spherical Minkowski coordinates $(m, r_M, \theta, \phi)$ reads as  

$$ds^2 = dm^2 - dr_M^2 - r_M^2 d\Omega^2.$$  

(3.1)  

2. Also the line element of $S^2$ is familiar:  

$$ds^2 = -R^2 (d\Theta^2 + \sin^2(\theta) d\Phi^2).$$  

(3.2)  

3. By visiting in Wikipedia one learns that in spherical coordinate the line element of AdS$_4$/dS$_4$/dS is given by  

$$ds^2 = A(r) dt^2 - \frac{1}{A(r)} dr^2 - r^2 d\Omega^2,$$  

$$A(r) = 1 + \epsilon y^2, \quad y = \frac{r}{r_0},$$  

$$\epsilon = 1 \text{ for AdS}_4, \quad \epsilon = -1 \text{ for dS}_4.$$  

(3.3)  

4. From these formulas it is easy to see that the ansatz is of the same general form as for the imbedding of Schwartzschild-Nordstöm metric:  

$$m = \Lambda t + h(y), \quad r_M = r,$$  

$$\Theta = s(y), \quad \Phi = \omega(t + f(y)).$$  

(3.4)  

The non-trivial conditions on the components of the induced metric are given by  

$$g_{tt} = \Lambda^2 - x^2 \sin^2(\Theta) = A(r),$$  

$$g_{tr} = \frac{1}{r_0} \left[ \frac{\Lambda}{y} - x^2 \sin^2(\theta) \frac{df}{dr} \right] = 0,$$  

$$g_{rr} = \frac{1}{r_0^2} \left[ \frac{df}{dy} - 1 - x^2 \sin^2(\theta) \frac{df}{dy} - R^2 \frac{d\Theta}{dy} \right] = - \frac{1}{A(r)},$$  

$$x = R \omega.$$  

(3.5)  

By some simple algebraic manipulations one can derive expressions for $\sin(\Theta)$, $df/dr$ and $dh/dr$.  

1. For $\Theta(r)$ the equation for $g_{tt}$ gives the expression  

$$\sin(\Theta) = \pm \frac{p^{1/2}}{x},$$  

$$P = \Lambda^2 - A = \Lambda^2 - 1 - \epsilon y^2.$$  

(3.6)  

The condition $0 \leq \sin^2(\Theta) \leq 1$ gives the conditions  

$$(\Lambda^2 - x^2 - 1)^{1/2} \leq y \leq (\Lambda^2 - 1)^{1/2} \quad \text{for } \epsilon = 1 \text{ (AdS}_4),$$  

$$(-\Lambda^2 + 1)^{1/2} \leq y \leq (x^2 + 1 - \Lambda^2)^{1/2} \quad \text{for } \epsilon = -1 \text{ (dS}_4).$$  

(3.7)  

Only a spherical shell is possible in both cases. The model for the final state of star considered in [K4] predicted similar layer layer like structure and inspired the proposal that stars quite generally have an onionlike structure with radii of various shells characterize by p-adic length scale hypothesis and thus coming in some powers of $\sqrt{2}$. This brings in mind also Titius-Bode law.
2. From the vanishing of $g_{tr}$ one obtains

$$\frac{dh}{dy} = \frac{P}{K} \frac{df}{dy} .$$

(3.8)

3. The condition for $g_{rr}$ gives

$$(\frac{df}{dy})^2 = \frac{r_0^2}{AP} \left[ A^{-1} - R^2 (\frac{d\Theta}{dy})^2 \right] .$$

(3.9)

Clearly, the right-hand side is positive if $P \geq 0$ holds true and $R d\Theta/dy$ is small. One can express $d\Theta/dy$ using chain rule as

$$(\frac{d\Theta}{dy})^2 = \frac{x^2 y^2}{P(P-x^2)} .$$

(3.10)

One obtains

$$(\frac{df}{dy})^2 = A r_0^2 \frac{y^2}{AP} \left[ \frac{1}{1+y^2} - x^2 (\frac{R}{r_0})^2 \frac{1}{P(P-x^2)} \right] .$$

(3.11)

The right hand side of this equation is non-negative for certain range of parameters and variable $y$. Note that for $r_0 \gg R$ the second term on the right hand side can be neglected. In this case it is easy to integrate $f(y)$.

The conclusion is that both AdS$_4$ and dS$_4$ allow a local imbedding as a vacuum extremal. Whether also an imbedding as a non-vacuum preferred extremal to $M^4 \times S^2$, $S^2$ a homologically non-trivial geodesic sphere is possible, is an interesting question.

4 Generalizing Ricci flow to Maxwell flow for 4-geometries and Kähler flow for space-time surfaces

The notion of Ricci flow has played a key part in the geometrization of topological invariants of Riemann manifolds. I certainly did not have this in mind when I choose to call my unification attempt "Topological Geometrodynamics" but this title strongly suggests that a suitable generalization of Ricci flow could play a key role in the understanding of also TGD.

4.1 Ricci flow and Maxwell flow for 4-geometries

The observation about constancy of 4-D curvature scalar for preferred extremals inspires a generalization of the well-known volume preserving Ricci flow [A4] introduced by Richard Hamilton. Ricci flow is defined in the space of Riemann metrics as

$$\frac{dg_{\alpha\beta}}{dt} = -2R_{\alpha\beta} + 2 \frac{R_{avg}}{D} g_{\alpha\beta} .$$

(4.1)

Here $R_{avg}$ denotes the average of the scalar curvature, and $D$ is the dimension of the Riemann manifold. The flow is volume preserving in average sense as one easily checks ($\langle (g^{\alpha\beta} g_{\alpha\beta})/dt \rangle = 0$). The volume preserving property of this flow allows to intuitively understand that the volume of a 3-manifold in the asymptotic metric defined by the Ricci flow is topological invariant. The fixed points of the flow serve
as canonical representatives for the topological equivalence classes of 3-manifolds. These 3-manifolds (for instance hyperbolic 3-manifolds with constant sectional curvatures) are highly symmetric. This is easy to understand since the flow is dissipative and destroys all details from the metric.

What happens in the recent case? The first thing to do is to consider what might be called Maxwell flow in the space of all 4-D Riemann manifolds allowing Maxwell field.

1. First of all, the vanishing of the trace of Maxwell’s energy momentum tensor codes for the volume preserving character of the flow defined as

\[ \frac{dg_{\alpha\beta}}{dt} = T_{\alpha\beta} \quad (4.2) \]

Taking covariant divergence on both sides and assuming that \( d/dt \) and \( D_\alpha \) commute, one obtains that \( T_{\alpha\beta} \) is divergenceless.

This is true if one assumes Einstein’s equations with cosmological term. This gives

\[ \frac{dg_{\alpha\beta}}{dt} = kG_{\alpha\beta} + \Lambda g_{\alpha\beta} = kR_{\alpha\beta} + (\frac{-kR}{2} + \Lambda)g_{\alpha\beta} \quad (4.3) \]

The trace of this equation gives that the curvature scalar is constant. Note that the value of the Kähler coupling strength plays a highly non-trivial role in these equations and it is quite possible that solutions exist only for some critical values of \( \alpha_K \). Quantum criticality should fix the allow value triplets \((G, \Lambda, \alpha_K)\) apart from overall scaling

\[ (G, \Lambda, \alpha_K) \rightarrow (xG, \Lambda/x, x\alpha_K) \]

Fixing the value of \( G \) fixes the values remaining parameters at critical points. The rescaling of the parameter \( t \) induces a scaling by \( x \).

2. By taking trace one obtains the already mentioned condition fixing the curvature to be constant, and one can write

\[ \frac{dg_{\alpha\beta}}{dt} = kR_{\alpha\beta} - \Lambda g_{\alpha\beta} \quad (4.4) \]

Note that in the recent case \( R_{\text{avg}} = R \) holds true since curvature scalar is constant. The fixed points of the flow would be Einstein manifolds \([A1, A6]\) satisfying

\[ R_{\alpha\beta} = \frac{\Lambda}{k} g_{\alpha\beta} \quad (4.5) \]

3. It is by no means obvious that continuous flow is possible. The condition that Einstein-Maxwell equations are satisfied might pick up from a completely general Maxwell flow a discrete subset as solutions of Einstein-Maxwell equations with a cosmological term. If so, one could assign to this subset a sequence of values \( t_n \) of the flow parameter \( t \).

4. I do not know whether 3-dimensionality is somehow absolutely essential for getting the topological classification of closed 3-manifolds using Ricci flow. This ignorance allows me to pose some innocent questions. Could one have a canonical representation of 4-geometries as spaces with constant Ricci scalar? Could one select one particular Einstein space in the class four-metrics and could the ratio \( \Lambda/k \) represent topological invariant if one normalizes metric or curvature scalar suitably. In the 3-dimensional case curvature scalar is normalized to unity. In the recent case this normalization would give \( k = 4\Lambda \) in turn giving \( R_{\alpha\beta} = g_{\alpha\beta}/4 \). Does this mean that there is only single fixed point in local sense, analogous to black hole toward which all geometries are driven by the Maxwell flow? Does this imply that only the 4-volume of the original space would serve as a topological invariant?
4.2 Maxwell flow for space-time surfaces

One can consider Maxwell flow for space-time surfaces too. In this case Kähler flow would be the appropriate term and provides families of preferred extremals. Since space-time surfaces inside CD are the basic physical objects in TGD framework, a possible interpretation of these families would be as flows describing physical dissipation as a four-dimensional phenomenon polishing details from the space-time surface interpreted as an analog of Bohr orbit.

1. The flow is now induced by a vector field $j^k(x,t)$ of the space-time surface having values in the tangent bundle of imbedding space $M^4 \times CP_2$. In the most general case one has Kähler flow without the Einstein equations. This flow would be defined in the space of all space-time surfaces or possibly in the space of all extremals. The flow equations reduce to

$$h_{kl} D_\alpha j^k(x,t) \partial_\beta h^l = \frac{1}{2} T_{\alpha\beta} ,$$

$$T^{\alpha\beta} = \frac{1}{16 \pi \alpha K} \left[ J^\alpha J^\beta - \frac{1}{4} J^\alpha J^\alpha J^\beta \right].$$

(4.6)

The left hand side is the projection of the covariant gradient $D_\alpha j^k(x,t)$ of the flow vector field $j^k(x,t)$ to the tangent space of the space-time surface. $D_\alpha$ is covariant derivative taking into account that $j^k$ is imbedding space vector field. For a fixed point space-time surface this projection must vanish assuming that this space-time surface reachable. A good guess for the asymptotia is that the divergence of Maxwell energy momentum tensor vanishes and that Einstein’s equations with cosmological constant are well-defined.

Asymptotes corresponds to vacuum extremals. In Euclidian regions $CP_2$ type vacuum extremals and in Minkowskian regions to any space-time surface in any 6-D sub-manifold $M^4 \times Y^2$, where $Y^2$ is Lagrangian sub-manifold of $CP_2$ having therefore vanishing induced Kähler form. Symplectic transformations of $CP_2$ combined with diffeomorphisms of $M^4$ give new Lagrangian manifolds. One would expect that vacuum extremals are approached but never reached at second extreme for the flow.

If one assumes Einstein’s equations with a cosmological term, allowed vacuum extremals must be Einstein manifolds. For $CP_2$ type vacuum extremals this is the case. It is quite possible that these fixed points do not actually exist in Minkowskian sector, and could be replaced with more complex asymptotic behavior such as limit, chaos, or strange attractor.

2. The flow could be also restricted to the space of preferred extremals. Assuming that Einstein Maxwell equations indeed hold true, the flow equations reduce to

$$h_{kl} D_\alpha j^k(x,t) \partial_\beta h^l = \frac{1}{2} (k R_{\alpha\beta} - \Lambda g_{\alpha\beta}) .$$

(4.7)

Preferred extremals would correspond to a fixed sub-manifold of the general flow in the space of all 4-surfaces.

3. One can also consider a situation in which $j^k(x,t)$ is replaced with $j^k(h,t)$ defining a flow in the entire imbedding space. This assumption is probably too restrictive. In this case the equations reduce to

$$(D_r j^l(x,t) + D_l j^r) \partial_\alpha h^r \partial_\beta h^l = k R_{\alpha\beta} - \Lambda g_{\alpha\beta} .$$

(4.8)

Here $D_r$ denotes covariant derivative. Asymptotia is achieved if the tensor $D_k j^l + D_l j^k$ becomes orthogonal to the space-time surface. Note for that Killing vector fields of $H$ the left hand side vanishes identically. Killing vector fields are indeed symmetries of also asymptotic states.
4.3 Dissipation, self organization, transition to chaos, and coupling constant evolution

It must be made clear that the existence of a continuous flow in the space of preferred extremals might be too strong a condition. Already the restriction of the general Maxwell flow in the space of metrics to solutions of Einstein-Maxwell equations with cosmological term might lead to discretization, and the assumption about representability as 4-surface in $M^4 \times CP_2^2$ would give a further condition reducing the number of solutions. On the other hand, one might consider a possibility of a continuous flow in the space of constant Ricci scalar metrics with a fixed 4-volume and having hyperbolic spaces as the most symmetric representative.

4.3 Dissipation, self organization, transition to chaos, and coupling constant evolution

A beautiful connection with concepts like dissipation, self-organization, transition to chaos, and coupling constant evolution suggests itself.

1. It is not at all clear whether the vacuum extremal limits of the preferred extremals can correspond to Einstein spaces except in special cases such as $CP_2^2$ type vacuum extremals isometric with $CP_2^2$. The imbeddability condition however defines a constraint force which might well force asymptotically more complex situations such as limit cycles and strange attractors. In ordinary dissipative dynamics an external energy feed is essential prerequisite for this kind of non-trivial self-organization patterns.

In the recent case the external energy feed could be replaced by the constraint forces due to the imbeddability condition. It is not too difficult to imagine that the flow (if it exists!) could define something analogous to a transition to chaos taking place in a stepwise manner for critical values of the parameter $t$. Alternatively, these discrete values could correspond to those values of $t$ for which the preferred extremal property holds true for a general Maxwell flow in the space of 4-metrics. Therefore the preferred extremals of Kähler action could emerge as one-parameter (possibly discrete) families describing dissipation and self-organization at the level of space-time dynamics.

2. For instance, one can consider the possibility that in some situations Einstein’s equations split into two mutually consistent equations of which only the first one is independent

\[
\begin{align*}
  xJ^\alpha_{\nu}J^{\nu\beta} &= R^{\alpha\beta}, \\
  L_K &= xJ^\alpha_{\nu}J^{\nu\beta} = 4\Lambda, \\
  x &= \frac{1}{16\pi\alpha_K}.
\end{align*}
\]

(4.9)

Note that the first equation indeed gives the second one by tracing. This happens for $CP_2^2$ type vacuum extremals.

Kähler action density would reduce to cosmological constant which should have a continuous spectrum if this happens always. A more plausible alternative is that this holds true only asymptotically. In this case the flow equation could not lead arbitrary near to vacuum extremal, and one can think of situation in which $L_K = 4\Lambda$ defines an analog of limiting cycle or perhaps even strange attractor. In any case, the assumption would allow to deduce the asymptotic value of the action density which is of utmost importance from calculational point of view: action would be simply $S_K = 4\Lambda V_4$ and one could also say that one has minimal surface with $\Lambda$ taking the role of string tension.

3. One of the key ideas of TGD is quantum criticality implying that Kähler coupling strength is analogous to critical temperature. Second key idea is that $p$-adic coupling constant evolution represents discretized version of continuous coupling constant evolution so that each $p$-adic prime would correspond a fixed point of ordinary coupling constant evolution in the sense that the 4-volume characterized by the $p$-adic length scale remains constant. The invariance of the geometric and thus geometric parameters of hyperbolic 4-manifold under the Kähler flow would conform with the interpretation as a flow preserving scale assignable to a given $p$-adic prime. The continuous evolution in question (if possible at all!) might correspond to a fixed $p$-adic
4.4 Does a 4-D counterpart of thermodynamics make sense?

The interpretation of the Kähler flow in terms of dissipation, the constancy of $R$, and almost constancy of $L_K$ suggest an interpretation in terms of 4-D variant of thermodynamics natural in zero energy ontology (ZEO), where physical states are analogs for pairs of initial and final states of quantum event are quantum superpositions of classical time evolutions. Quantum theory becomes a “square root” of thermodynamics so that 4-D analog of thermodynamics might even replace ordinary thermodynamics as a fundamental description. If so this 4-D thermodynamics should be qualitatively consistent with the ordinary 3-D thermodynamics.

1. The first naive guess would be the interpretation of the action density $L_K$ as an analog of energy density $e = E/V_3$ and that of $R$ as the analog to entropy density $s = S/V_3$. The asymptotic states would be analogs of thermodynamical equilibria having constant values of $L_K$ and $R$.

2. Apart from an overall sign factor $\epsilon$ to be discussed, the analog of the first law $de = Tds - pdV/V$ would be

$$dL_K = kdR + \Lambda \frac{dV_4}{V_4}.$$ 

One would have the correspondences $S \rightarrow \epsilon RV_4$, $e \rightarrow \epsilon L_K$ and $k \rightarrow T$, $p \rightarrow -\Lambda$. $k \propto 1/G$ indeed appears formally in the role of temperature in Einstein’s action defining a formal partition function via its exponent. The analog of second law would state the increase of the magnitude of $\epsilon RV_4$ during the Kähler flow.

3. One must be very careful with the signs and discuss Euclidian and Minkowskian regions separately. Concerning purely thermodynamic aspects at the level of vacuum functional Euclidian regions are those which matter.

(a) For $CP_2$ type vacuum extremals $L_K \propto E^2 + B^2$, $R = \Lambda/k$, and $\Lambda$ are positive. In thermodynamical analogy for $\epsilon = 1$ this would mean that pressure is negative.

(b) In Minkowskian regions the value of $R = \Lambda/k$ is negative for $\Lambda < 0$ suggested by the large abundance of 4-manifolds allowing hyperbolic metric and also by cosmological considerations. The asymptotic formula $L_K = 4\Lambda$ considered above suggests that also Kähler action is negative in Minkowskian regions for magnetic flux tubes dominating in TGD inspired cosmology: the reason is that the magnetic contribution to the action density $L_K \propto E^2 - B^2$ dominates.

Consider now in more detail the 4-D thermodynamics interpretation in Euclidian and Minkowskian regions assuming that the the evolution by quantum jumps has Kähler flow as a space-time correlate.

1. In Euclidian regions the choice $\epsilon = 1$ seems to be more reasonable one. In Euclidian regions $-\Lambda$ as the analog of pressure would be negative, and asymptotically (that is for $CP_2$ type vacuum extremals) its value would be proportional to $\Lambda \propto 1/GR^2$, where $R$ denotes $CP_2$ radius defined by the length of its geodesic circle.
5. Could correlation functions, S-matrix, and coupling constant evolution be coded the statistical properties of preferred extremals?

A possible interpretation for negative pressure is in terms of string tension effectively inducing negative pressure (note that the solutions of the modified Dirac equation indeed assign a string to the wormhole contact). The analog of the second law would require the increase of $RV_4$ in quantum jumps. The magnitudes of $L_K$, $R$, $V_4$ and $\Lambda$ would be reduced and approach their asymptotic values. In particular, $V_4$ would approach asymptotically the volume of $CP_2$.

2. In Minkowskian regions Kähler action contributes to the vacuum functional a phase factor analogous to an imaginary exponent of action serving in the role of Morse function so that thermodynamics interpretation can be questioned. Despite this one can check whether thermodynamic interpretation can be considered. The choice $\epsilon = -1$ seems to be the correct choice now. $-\Lambda$ would be analogous to a negative pressure whose gradually decreases. In 3-D thermodynamics it is natural to assign negative pressure to the magnetic flux tube like structures as their effective string tension defined by the density of magnetic energy per unit length. $-R \geq 0$ would entropy and $-L_K \geq 0$ would be the analog of energy density.

Assume the recent view about state function reduction explaining how the arrow of geometric time is induced by the quantum jump sequence defining experienced time [K1]. According to this view zero energy states are quantum superpositions over $CD$s of various size scales but with common tip, which can correspond to either the upper or lower light-like boundary of $CD$. The sequence of quantum jumps the gradual increase of the average size of $CD$ in the quantum superposition and therefore that of average value of $V_4$. On the other hand, a gradual decrease of both $-L_K$ and $-R$ looks physically very natural. If Kähler flow describes the effect of dissipation by quantum jumps in ZEO then the space-time surfaces would gradually approach nearly vacuum extremals with constant value of entropy density $-R$ but gradually increasing 4-volume so that the analog of second law stating the increase of $-RV_4$ would hold true.

3. The interpretation of $-R > 0$ as negentropy density assignable to entanglement is also possible and is consistent with the interpretation in terms of second law. This interpretation would only change the sign factor $\epsilon$ in the proposed formula. Otherwise the above arguments would remain as such.

5. Could correlation functions, S-matrix, and coupling constant evolution be coded the statistical properties of preferred extremals?

Quantum classical correspondence states that all aspects of quantum states should have correlates in the geometry of preferred extremals. In particular, various elementary particle propagators should have a representation as properties of preferred extremals. This would allow to realize the old dream about being able to say something interesting about coupling constant evolution although it is not yet possible to calculate the M-matrices and U-matrix. Hitherto everything that has been said about coupling constant evolution has been rather speculative arguments except for the general vision that it reduces to a discrete evolution defined by p-adic length scales. General first principle definitions are however much more valuable than ad hoc guesses even if the latter give rise to explicit formulas.

In quantum TGD and also at its QFT limit various correlation functions in given quantum state should code for its properties. By quantum classical correspondence these correlation functions should have counterparts in the geometry of preferred extremals. Even more: these classical counterparts for a given preferred extremal ought to be identical with the quantum correlation functions for the superposition of preferred extremals. This correspondence could be called quantum ergodicity by its analogy with ordinary ergodicity stating that the member of ensemble becomes representative of ensemble.

1. The marvelous implication of quantum ergodicity would be that one could calculate everything solely classically using the classical intuition - the only intuition that we have. Quantum ergodicity would also solve the paradox raised by the quantum classical correspondence for momentum.
5. Could correlation functions, S-matrix, and coupling constant evolution be coded the statistical properties of preferred extremals? 11
eigenstates. Any preferred extremal in their superposition defining momentum eigenstate should code for the momentum characterizing the superposition itself. This is indeed possible if every extremal in the superposition codes the momentum to the properties of classical correlation functions which are identical for all of them.

2. The only manner to possibly achieve quantum ergodicity is in terms of the statistical properties of the preferred extremals. It should be possible to generalize the ergodic theorem stating that the properties of statistical ensemble are represented by single space-time evolution in the ensemble of time evolutions. Quantum superposition of classical worlds would effectively reduce to single classical world as far as classical correlation functions are considered. The notion of finite measurement resolution suggests that one must state this more precisely by adding that classical correlation functions are calculated in a given UV and IR resolutions meaning UV cutoff defined by the smallest CD and IR cutoff defined by the largest CD present.

3. The skeptic inside me immediately argues that TGD Universe is 4-D spin glass so that this quantum ergodic theorem must be broken. In the case of the ordinary spin classes one has not only statistical average for a fixed Hamiltonian but a statistical average over Hamiltonians. There is a probability distribution over the coupling parameters appearing in the Hamiltonian. Maybe the quantum counterpart of this is needed to predict the physically measurable correlation functions.

Could this average be an ordinary classical statistical average over quantum states with different classical correlation functions? This kind of average is indeed taken in density matrix formalism. One could assign these contributions to different "phases".

4. Quantum classical correspondence in statistical sense would be very much like holography (now individual classical state represents the entire quantum state). Quantum ergodicity would pose a rather strong constraint on quantum states. This symmetry principle could actually fix the spectrum of zero energy states to a high degree and fix therefore the M-matrices given by the product of hermitian square root of density matrix and unitary S-matrix and unitary U-matrix having M-matrices as its orthonormal rows.

5. In TGD inspired theory of consciousness the counterpart of quantum ergodicity is the postulate that the space-time geometry provides a symbolic representation for the quantum states and also for the contents of consciousness assignable to quantum jumps between quantum states. Quantum ergodicity would realize this strongly self-referential looking condition. The positive and negative energy parts of zero energy state would be analogous to the initial and final states of quantum jump and the classical correlation functions would code for the contents of consciousness like written formulas code for the thoughts of mathematician and provide a sensory feedback.

How classical correlation functions should be defined?

1. General Coordinate Invariance and Lorentz invariance are the basic constraints on the definition. These are achieved for the space-time regions with Minkowskian signature and 4-D $M^4$ projection if linear Minkowski coordinates are used. This is equivalent with the contraction of the indices of tensor fields with the space-time projections of $M^4$ Killing vector fields representing translations. Accepting this generalization, there is no need to restrict oneself to 4-D $M^4$ projection and one can also consider also Euclidian regions identifiable as lines of generalized Feynman diagrams.

Quantum ergodicity very probably however forces to restrict the consideration to Minkowskian and Euclidian space-time regions and various phases associated with them. Also $CP_2$ Killing vector fields can be projected to space-time surface and give a representation for classical gluon fields. These in turn can be contracted with $M^4$ Killing vectors giving rise to gluon fields as analogs of graviton fields but with second polarization index replaced with color index.

2. The standard definition for the correlation functions associated with classical time evolution is the appropriate starting point. The correlation function $G_{XY}(\tau)$ for two dynamical variables
X(t) and Y(t) is defined as the average \( G_{XY}(\tau) = \frac{1}{T} \int_{T} X(t)Y(t+\tau)dt \) over an interval of length \( T \), and one can also consider the limit \( T \to \infty \). In the recent case one would replace \( \tau \) with the difference \( m_1 - m_2 = m \) of \( M^4 \) coordinates of two points at the preferred extremal and integrate over the points of the extremal to get the average. The finite time interval \( T \) is replaced with the volume of causal diamond in a given length scale. Zero energy state with given quantum numbers for positive and negative energy parts of the state defines the initial and final states between which the fields appearing in the correlation functions are defined.

3. What correlation functions should be considered? Certainly one could calculate correlation functions for the induced spinor connection given electro-weak propagators and correlation functions for \( CP^2 \) Killing vector fields giving correlation functions for gluon fields using the description in terms of Killing vector fields. If one can uniquely separate from the Fourier transform uniquely a term of form \( Z/(p^2 - m^2) \) by its momentum dependence, the coefficient \( Z \) can be identified as coupling constant squared for the corresponding gauge potential component and one can in principle deduce coupling constant evolution purely classically. One can imagine of calculating spinorial propagators for string world sheets in the same manner. Note that also the dependence on color quantum numbers would be present so that in principle all that is needed could be calculated for a single preferred extremal without the need to construct QFT limit and to introduce color quantum numbers of fermions as spin like quantum numbers (color quantum numbers corresponds to \( CP^2 \) partial wave for the tip of the CD assigned with the particle).

4. What about Higgs field? TGD in principle allows scalar and pseudo-scalars which could be called Higgs like states. These states are however not necessary for particle massivation although they can represent particle massivation and must do so if one assumes that QFT limit exist. p-Adic thermodynamics however describes particle massivation microscopically.

The problem is that Higgs like field does not seem to have any obvious space-time correlate. The trace of the second fundamental form is the obvious candidate but vanishes for preferred extremals which are both minimal surfaces and solutions of Einstein Maxwell equations with cosmological constant. If the string world sheets at which all spinor components except right handed neutrino are localized for the general solution ansatz of the modified Dirac equation, the corresponding second fundamental form at the level of imbedding space defines a candidate for classical Higgs field. A natural expectation is that string world sheets are minimal surfaces of space-time surface. In general they are however not minimal surfaces of the imbedding space so that one might achieve a microscopic definition of classical Higgs field and its vacuum expectation value as an average of one point correlation function over the string world sheet.

Many detailed speculations about coupling constant evolution to be discussed in the sections below must be taken as innovative guesses doomed to have the eventual fate of guesses. The notion of quantum ergodicity could however be one of the really deep ideas about coupling constant evolution comparable to the notion of p-adic coupling constant evolution. Quantum Ergodicity (briefly QE) would also state something extremely non-trivial also about the construction of correlation functions and S-matrix. Because this principle is so new, the rest of the chapter does not yet contain any applications of QE. This should not lead the reader to under-estimate the potential power of QE.

Mathematics

Books related to TGD


